

Variation Operator

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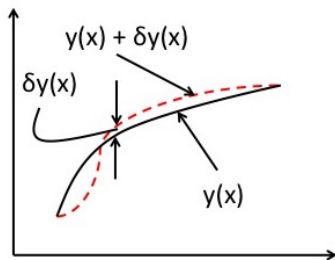
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Euler-Lagrange Equations Using Delta Operator

Consider a Function $y(x)$ and a Neighborhood function $\tilde{y}(x)$ to the function $y(x)$

$$\tilde{y}(x) = y(x) + \delta y(x)$$

$\delta \Rightarrow$ Delta Operator or Variation Operator



Euler-Lagrange Equations Using Delta Operator

$\delta y(x)$ = An infinitesimal, slowly varying, change to the function at a given x . It vanishes at those points where $y(x)$ is specified.

$\delta y(x)$ is not same as $\frac{dy}{dx}$, it is $\delta\left(\frac{dy}{dx}\right) = \frac{d}{dx}(\delta y)$. It has the following properties:

- $\delta(F \pm G) = \delta F \pm \delta G$
- $\delta(FG) = F\delta G + G\delta F$
- $\delta\left(\frac{F}{G}\right) = \frac{G\delta F - F\delta G}{G^2}$

Euler-Lagrange Equations Using Delta Operator

Consider a Functional

$$I = \int_{x_1}^{x_2} F(x, y, y') dx$$

Here $y = y(x)$ (assumed to be continuous in $x_1 < x < x_2$), and y' indicates derivative of y with respect to x ; and F , called the Lagrange function or Lagrangian, is a function of x , y and y' .

Objective: Determine $y(x)$ that will make I to be stationary.

$\delta y(x)$ vanishes at the boundary where the function $y(x)$ (essential boundary condition) is specified.

Euler-Lagrange Equations (contd...)

Variation of I

$$\begin{aligned}\delta I &= \int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' \right) dx + \text{higher order terms} \\ &= \delta^{(1)} I + \text{higher order terms}\end{aligned}$$

For I to have a stationary value, the first term on the right hand side, called the first variation of I , represented as $\delta^{(1)}$ must vanish (i.e $\delta^{(1)} = 0$). This yields:

$$\int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' \right) dx = 0$$

Euler-Lagrange Equations (contd...)

Integrating the second term by parts:

$$\int_{x_1}^{x_2} \frac{\partial F}{\partial y'} \delta y'(x) dx = \left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \delta y(x) dx$$

Using the following Property of delta operator

$$\delta \left(\frac{dy}{dx} \right) = \frac{d(\delta y(x))}{dx}$$

Stationary condition for I $\delta^{(1)} = 0$

$$\left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_1}^{x_2} + \int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] \delta y(x) dx = 0$$

Euler-Lagrange Equations (contd...)

$\delta y(x_1)$ and $\delta y(x_2)$ are arbitrary and independent

Above equation must be true for all values of $\delta y(x_1)$ and $\delta y(x_2)$ including when both are zero. This means the second term must be zero by itself.

$$\left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_1}^{x_2} = 0$$

$$\int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] \delta y(x) dx = 0$$

Since $\delta y(x)$ is an arbitrary and a slowly varying function of x

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$$

Euler-Lagrange Equations (contd...)

Euler-Lagrange equation for the given Lagrangian, F .

$\delta y(x_1)$ and $\delta y(x_2)$ are independent from each other and are completely arbitrary

$$\left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_1} = 0$$

$$\left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_2} = 0$$

Either $\delta y(x_1) = 0$ ($y(x_1)$ specified), or $\frac{\partial F}{\partial y'} = 0$ at x_1 ,

and either $\delta y(x_2) = 0$ ($y(x_2)$ specified), or $\frac{\partial F}{\partial y'} = 0$ at x_2 .

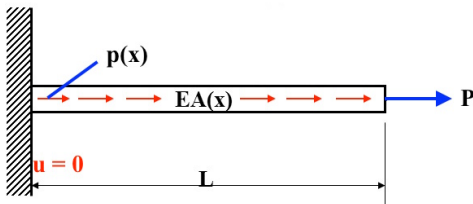
Euler-Lagrange Equations (contd...)

At a given end, $\frac{\partial F}{\partial y'}$ may be specified. For such cases the functional I will have additional term $G(y_1)$ or $(G(y_2))$.

For some cases, When $y(x)$ is specified at an end, we term it to be an **Essential Boundary Condition**

When $\frac{\partial F}{\partial y'} = 0$, it is called **Natural Boundary Condition**

Example



Total potential energy (Π)

$$\Pi = U + V$$

U = Strain Energy of the bar

$$= \frac{1}{2} \int_0^L [EA(x)] \left(\frac{du}{dx} \right)^2 dx$$

V = Potential of the applied load

$$= - \int_0^L p(x) u(x) dx - Pu(x = L)$$

Example (contd...)

Principle of minimum total potential energy

$\delta^{(1)}\Pi = 0$, $\delta^{(1)}\Pi$ is the first variation of the total potential energy

$$\begin{aligned}\delta^{(1)}\Pi = & \frac{1}{2} \int_0^L [EA(x)] \cdot 2 \cdot \left(\frac{du}{dx} \right) \delta \left(\frac{du}{dx} \right) dx \\ & - \int_0^L p(x) \delta u(x) dx - P \delta u(x=L)\end{aligned}$$

By using integration by parts, and the fact that

$$\delta \left(\frac{du}{dx} \right) = \frac{d(\delta u)}{dx}$$

Example (contd...)

The first variation of Π becomes:

$$\begin{aligned}\delta^{(1)}\Pi &= EA(x) \left(\frac{du}{dx} \right) \delta u \Big|_0^L - \int_0^L \frac{d}{dx} \left[EA(x) \left(\frac{du}{dx} \right) \right] \delta u(x) dx \\ &\quad - \int_0^L p(x) \delta u(x) dx - P \delta u(x=L)\end{aligned}$$

At $x = 0$, the displacement u is specified (an essential boundary condition). This renders $\delta u = 0$ at $x = 0$.

Example (contd...)

For Π to be stationary,

$$\left[EA(x) \left(\frac{du}{dx} \right) \Big|_{x=L} - P \right] \delta u \Big|_{x=L} - \int_0^L \left[\frac{d}{dx} \left\{ EA(x) \left(\frac{du}{dx} \right) \right\} + p(x) \right] \delta u(x) dx = 0$$

The above equation must be satisfied for all values of $\delta u(x = L)$, including $\delta u(x = L) = 0$

Euler Lagrange Equation

$$\frac{d}{dx} \left[EA(x) \left(\frac{du}{dx} \right) \right] + p(x) = 0$$

The associated boundary condition at $x = L$ is given as:

$$EA(x) \left(\frac{du}{dx} \right) = P$$

The boundary condition at $x = L$ is a **natural boundary condition**.

Lagrangian with Second Derivatives

Euler-Bernoulli Beam Strain energy contains the square of the second derivative of the transverse deflection

Euler-Lagrange equation will be a fourth-order ordinary differential equation with two boundary conditions at each end.

Consider

$$I = \int_{x_1}^{x_2} F(x, y, y', y'') dx$$

Here $y = y(x)$ (assumed to be continuous in $x_1 < x < x_2$), y' indicates derivative of $y(x)$ with respect to x ; y'' indicates second derivative of $y(x)$ with respect to x , and F , called the Lagrange function or Lagrangian, is a function of x, y, y' and y''

Objective:

Determine $y(x)$ that will make I to be stationary.

Lagrangian with Second Derivatives (contd...)

$$\begin{aligned}\delta I &= \int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' + \frac{\partial F}{\partial y''} \delta y'' \right) dx \\ &\quad + \text{higher order terms} \\ &= \delta^{(1)} I + \text{higher order terms}\end{aligned}$$

For I to have a stationary value $\delta^{(1)}$ must vanish i.e. $\delta^{(1)} = 0$.

$$\int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' + \frac{\partial F}{\partial y''} \delta y'' \right) dx = 0$$

Lagrangian with Second Derivatives (contd...)

Integrating the second term by parts and using the fact that

$$\delta \left(\frac{dy}{dx} \right) = \frac{d(\delta y)}{dx}$$

yield:

$$\int_{x_1}^{x_2} \frac{\partial F}{\partial y'} \delta y'(x) dx = \left. \frac{\partial F}{\partial y'} \delta y(x) \right|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \delta y(x) dx$$

Similarly, integrating the third term by parts and using the fact that

$$\delta \left(\frac{d^2 y}{dx^2} \right) = \frac{d^2(\delta y)}{dx^2}$$

yield:

$$\int_{x_1}^{x_2} \frac{\partial F}{\partial y''} \delta y''(x) dx = \left. \frac{\partial F}{\partial y''} \delta y'(x) \right|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \delta y'(x) dx$$

Lagrangian with Second Derivatives (contd...)

Performing the integration by parts one more time, we get:

$$\begin{aligned}\int_{x_1}^{x_2} \frac{\partial F}{\partial y''} \delta y''(x) dx &= \left. \frac{\partial F}{\partial y''} \delta y'(x) \right|_{x_1}^{x_2} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \delta y(x) \Big|_{x_1}^{x_2} \\ &\quad + \int_{x_1}^{x_2} \frac{d^2}{dx^2} \left(\frac{\partial F}{\partial y''} \right) \delta y(x) dx\end{aligned}$$

Combining:

$$\begin{aligned}\delta^{(1)} I &= \left. \frac{\partial F}{\partial y''} \delta y'(x) \right|_{x_1}^{x_2} + \left[\frac{\partial F}{\partial y'} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) \Big|_{x_1}^{x_2} \\ &\quad + \int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) + \frac{d^2}{dx^2} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) dx\end{aligned}$$

Lagrangian with Second Derivatives (contd...)

For I to be stationary, we want $\delta^{(1)}I = 0$. This yields:

$$\begin{aligned} \frac{\partial F}{\partial y''} \delta y'(x) \Big|_{x_1}^{x_2} + \left[\frac{\partial F}{\partial y'} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) \Big|_{x_1}^{x_2} \\ + \int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) + \frac{d^2}{dx^2} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) dx = 0 \end{aligned}$$

Note that $\delta y(x)$, and $\delta y'(x)$ at a given end are arbitrary and independent of each other.

The above equation has to be satisfied for all possible values of $\delta y(x)$, and $\delta y'(x)$ at the two ends, including all four of them being zero *i.e.* $\delta y(x_1) = 0$, $\delta y(x_2) = 0$, $\delta y'(x_1) = 0$, and $\delta y'(x_2) = 0$.

Lagrangian with Second Derivatives (contd...)

This means:

$$\frac{\partial F}{\partial y''} \delta y'(x) \Big|_{x_1}^{x_2} = 0$$

$$\left[\frac{\partial F}{\partial y'} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) \Big|_{x_1}^{x_2} = 0$$

$$\int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) + \frac{d^2}{dx^2} \left(\frac{\partial F}{\partial y''} \right) \right] \delta y(x) dx = 0$$

Lagrangian with Second Derivatives (contd...)

From the First Equation, either

$$\delta y' = 0 \quad \text{or} \quad \frac{\partial F}{\partial y''} = 0$$

First Condition

y is specified \rightarrow Essential Boundary Condition

Second Condition \rightarrow Natural or Force Boundary Condition

General Case

$\frac{\partial F}{\partial y''}$ may be specified and additional (boundary) term y' in the function

Lagrangian with Second Derivatives (contd...)

From the Second Equation, either

$$\delta y = 0 \quad \text{or} \quad \left[\frac{\partial F}{\partial y'} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \right] = 0$$

First Condition

y is specified \rightarrow Essential Boundary Condition

Second Condition \rightarrow Natural or Force Boundary Condition

General Case

$\left[\frac{\partial F}{\partial y'} - \frac{d}{dx} \left(\frac{\partial F}{\partial y''} \right) \right]$ may be specified and additional (boundary) term y in the function

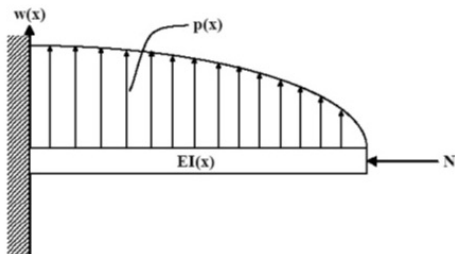
Lagrangian with Second Derivatives (contd...)

From the Third Term also being zero, the Euler-Lagrange equation is obtained by using the fact that variation $\delta y(x)$ is a slowly varying, arbitrary function of x .

$$\left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) + \frac{d^2}{dx^2} \left(\frac{\partial F}{\partial y''} \right) \right] = 0$$

The Euler-Lagrange equation for this case will be a fourth-order ordinary differential equation.

Example



The Total Potential energy of the system is:

$$\Pi = \frac{1}{2} \int_0^L \left[EI(x) \left(\frac{d^2 w}{dx^2} \right)^2 - N \left(\frac{dw}{dx} \right)^2 \right] dx - \int_0^L p(x) w(x) dx$$

Determine the Euler-Lagrange equations along with the associated boundary conditions.

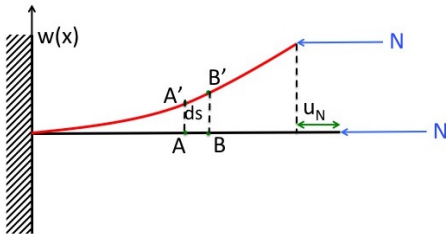
Example (contd...)

The **Strain Energy** of the system is,

$$U = \frac{1}{2} \int_0^L \left[EI(x) \left(\frac{d^2 w}{dx^2} \right)^2 \right] dx$$

The Potential of Compressive Load is,

Let u_N be the inward axial displacement of the right end due to transverse deflection.

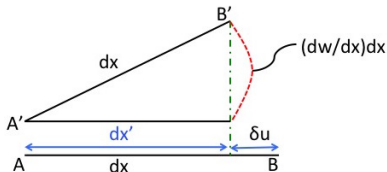


Example (contd...)

Now in order to determine u_N , zooming into the ' AB ' section of the beam,

Here,

$$\delta u = dx - dx'$$



From the figure,

$$\begin{aligned} &= \left[dx - \sqrt{dx^2 - \left(\frac{dw}{dx} \right)^2 dx^2} \right] \\ &\simeq dx \left[1 - 1 + \frac{1}{2} \left(\frac{dw}{dx} \right)^2 \right] \simeq \frac{1}{2} \left(\frac{dw}{dx} \right)^2 dx \\ u_N &= \frac{1}{2} \int_0^L \left(\frac{dw}{dx} \right)^2 dx \end{aligned}$$

Example (contd...)

We have

$$F = \frac{1}{2} \left[EI(x) \left(\frac{d^2 w}{dx^2} \right)^2 - N \left(\frac{dw}{dx} \right)^2 \right] - p(x)w(x)$$

$$\frac{dF}{dw} = -p$$

$$\frac{dF}{dw'} = \frac{1}{2} \cdot 2(-N) \frac{dw}{dx}$$

$$\frac{dF}{dw''} = \frac{1}{2} \cdot 2EI(x) \frac{d^2 w}{dx^2}$$

Example

The governing equation thus becomes:

$$-p(x) - \frac{d}{dx} \left[-N \frac{dw}{dx} \right] + \frac{d^2}{dx^2} \left[EI(x) \frac{d^2 w}{dx^2} \right] = 0$$

This simplifies to:

$$\frac{d^2}{dx^2} \left[EI(x) \frac{d^2 w}{dx^2} \right] + \frac{d}{dx} \left[N \frac{dw}{dx} \right] = p(x)$$

Example (contd...)

The boundary conditions at $x = L$ are:

$$N \frac{dw}{dx} + \frac{d}{dx} \left[EI(x) \frac{d^2 w}{dx^2} \right] = 0$$

$$EI(x) \frac{d^2 w}{dx^2} = 0$$

Both the boundary conditions at $x = L$ are **natural boundary conditions**.

The boundary conditions at $x = 0$ are:

$$\begin{aligned} w &= 0 \\ \frac{dw}{dx} &= 0 \end{aligned}$$

The two boundary conditions at $x = 0$ are both **essential boundary conditions**.